HOMEOSTASIS IN CHEMICAL REACTION PATHWAYS

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ABSTRACT. We consider stochastic models of chemical reaction networks with time dependent input rates and several types of molecules. We prove that, in despite of strong time dependence of input rates, there is a kind of homeostasis phenomenon: far away from input nodes the mean numbers of molecules of each type become approximately constant (do not depend on time).

Key words and phrases: stochastic models of biological systems, almost periodic input rates, stochastic networks, thermodynamic limit

1. Introduction and the Model

The problem has arised from thinking over homeostasis phenomenon. Many biological phenomena can be described in terms of the concentrations or the mean numbers $m_{\alpha,v}$ of molecules of type v in different parts (compartments) α of the biological organism. Compartments can be imagined as space separated parts of the cell or biological organism, where some system of chemical reactions takes place. As the result of chemical reactions in a given compartment, molecules can change their types and leave the compartment to arrive at another one. One of many features of the homeostasis is that large fluctuations of external media become smoothed inside of the organism

Consider a directed graph G with $M=N_1J$ vertices, each vertex is a pair $i=(\alpha,v)$, where v=1,...,J represent molecule types, and $\alpha=1,...,N_1$ are the compartments. We will define a continuous time Markov or, more generally, non Markov random process, the states of which are the vectors $(\xi_{\alpha,v},\alpha=1,...,N_1;v=1,...,J)$, $\xi_{\alpha,v}\in Z_+$. This process models chemical reaction networks, where the number $\xi_{\alpha,v}$ is interpreted as the number of molecules of type v in the compartment α . Edges $((\alpha,v),(\alpha,v'))$ of the network correspond to some mechanism of the chemical reaction in the compartment α , edges $((\alpha,v),(\alpha',v))$ correspond to molecule transportation from the compartment α to the compartment α' .

This random process is defined by the following transition rates:

- (1) Molecules of type v arrive from exterior with (time dependent) input rates $\lambda_i = \lambda_{\alpha,v}(t)$ to the compartment α . If λ_i is not identically zero then the vertex is called input vertex. Input rates are assumed to be almost periodic (in particular, periodic) or stationary stochastic;
- (2) In Markov case any molecule at the vertex i with rate $\lambda_{ij} = \lambda_{(\alpha,v)(\beta,w)}$, jumps to the vertex j, and with rate $\lambda_{i,0}$ leaves the network;

(3) In non Markov case another language is more convenient. After a molecule appears at vertex i, it decides immediately with (time independent) probabilities $p_{ij} = p_{(\alpha,v)(\beta,w)}$, to go from the compartment α to the compartment β , simultaneously changing its type v to w. Another possibility for this molecule is to leave the system with the probability p_{i0} . Obviously it should be

$$\sum_{i} p_{ij} + p_{i0} = 1$$

We assume that G has (one) directed edge from i to j iff p_{ij} is not identically zero. After this decision, some random time τ_{ij} is necessary until the molecule from the node i appears in the node j after this time. All these random times are mutually independent and have distribution functions $F_{i,j}$, which depend only on i,j. Note that in the Markov case $F_{i,j}$ are exponential with parameter λ_{ij} .

Now we explain our main result intuitively together with its biological interpretation. Our goal is to discuss homeostasis phenomenon in such networks. General mathematical notion of the homeostasis, as one of the most fundamental biological concepts, may not exist at all. However, inside formal models one may suggest definitions of the homeostasis, which could be reasonable for a given class of models. We think moreover that the notion of the homeostasis we suggest here could concern other applied networks.

In most of the paper, all rates, except input rates, are assumed to be constant (time independent). There are many reasons for such assumption. For example, the rates of unary reactions normally depend on the concentrations of the enzymes, catalyzing the corresponding unimolecular reaction. Enzyme concentrations are governed by a process of enzyme synthesis (DNA-RNA-ribosome-protein) which has slower time scale. That is why on the properly chosen time scale these rates are approximately constant. Moreover, in the vicinity of thermodynamic equilibrium (minimal Gibbs free energy), the reaction rates weakly depend on time, see [3].

Roughly speaking, we prove that (in despite of strong dependence of input rates, and thus of the concentrations in the input nodes, on time) there is a kind of homeostasis phenomenon: far away from input nodes the means

$$m_i(t) = E\xi_{\alpha,v}(t)$$

become approximately constant (do not depend on time). In real biology there exist, surely, more complicated control mechanisms. The goal of this paper is to show that even such natural rough mechanism like randomness sometimes effectively controls homeostasis.

2. Almost Periodic Input

A path of length n from vertex i to vertex j is a sequence $e_1 = (i_1, j_1), ..., e_n = (i_n, j_n)$ of directed edges such that $i_1 = i, j_n = j$ and $j_k = i_{k+1}$ for any k = 1, ..., n-1. Denote $\mathbf{L}_n(i,j)$ the set of all paths from i to j of length n.

We introduce a non negative measure P_n on the set $\mathbf{L}_n(i,j)$. For any path $L = ((i_1, i_2), ..., (i_{n-1}, i_n)), i_1 = i, i_n = j$, this measure is

$$(2.1) P_n(L) = p_{i_1 i_2} ... p_{i_{n-1} i_n}$$

Then

$$p^{(n)}(i,j) = \sum_{L \in \mathbf{L}_n(i,j)} P_n(L)$$

is the *n*-step transition probability for the discrete time (improper) Markov chain on the set of vertices of the graph G with the transition probability matrix (p_{ij}) .

It follows from the definitions, that each molecule, independently of others, performs random walk $\eta(t)$ on the graph, where $\eta(t)$ is a semi-Markov process. The distribution of this process is uniquely defined by transition probabilities p_{ij} and distribution functions F_{ij} . Let

$$P_{ij}(t) = P(\eta(t) = j | \eta(0) = i)$$

Here we suppose that at the initial moment the sojourn time at state i is equal to 0 for all states i. Then,

(2.2)
$$P_{ij}(t) = \sum_{n=1}^{\infty} \sum_{L \in \mathbf{L}_n(i,j)} P_n(L) \int_0^s dF_L(u) (1 - F_j(t - u)),$$

where F_L is the convolution of distributions $F_{i_1,i_2},...,F_{i_{n-1},i_n}$ and

$$F_i(s) = \sum_{i} F_{ij}(s) p_{ij}$$

We will need some technical assumptions:

(1) In the first theorem we assume that $\lambda_i(t), -\infty < t < \infty$, are non negative uniform almost periodic functions in the sense of [10]. One of the equivalent definitions of uniform almost periodic functions is that there exists the following limit for all real σ :

$$c(\sigma, i) = \lim_{T \to \infty} \frac{1}{T} \int_0^T \lambda_i(s) \exp(-i \sigma s) ds$$

and, moreover, the function $c(\sigma, i)$ is different from 0 at most on a countable set of values σ . Let $\sigma_{k,i}$ be the values of σ , such that $c(\sigma_{k,i}, i) \neq 0$. Put $c_k(i) = c(\sigma_{k,i}, i)$. Note that $c(0, i) \neq 0$, as $\lambda_i(s)$ is non negative, and we denote $c(0, i) = c_0(i)$. We will assume that $\sum_{k=0}^{\infty} |c_k(i)| < \infty$. Then the Fourier series

(2.3)
$$\lambda_i(t) = \sum_{k=0}^{\infty} c_k(i) \exp(i \sigma_{k,i} t),$$

converges to the function $\lambda_i(t)$ uniformly over $-\infty < t < \infty$.

(2) There exist the means

$$\int_0^\infty s dF_{ij}(s) < \infty$$

(3) Assume

(2.4)
$$b(i,j) = \sum_{n>0} p^{(n)}(i,j) < \infty$$

For this it is necessary that the probability to leave the network $p_{i0} > 0$ at least for one i.

Let I be the set of input nodes. Define the following constants

$$d_{j} = \sum_{i \in I} \overline{\lambda_{i}} b(i, j) \mu_{j}, \overline{\lambda_{i}} = c_{0}(i), \mu_{j} = \int_{0}^{\infty} s dF_{i}(s)$$

Introduce the class $\mathcal{M}(C, a, \delta)$, where C, a > 0 and $0 < \delta < 1$, consisting of the networks which satisfy the above assumptions (1)-(3) and also the following conditions:

(i) the following inequality holds

$$\sum_{i \in I} \sum_{k=0}^{\infty} |c_k(i)| \le C;$$

- (ii) if $|\sigma| < a$ and $\sigma \neq 0$, then $c(\sigma, i) = 0$ for all i;
- (iii) for any a > 0 there is $0 < \delta < 1$ such that if $|\sigma| \ge a$ then $|\psi_{i,j}(\sigma)| \le \delta < 1$, where $\psi_{i,j}(\sigma)$ are characteristic functions of distributions F_{ij} ;

Theorem 1. Fix some network satisfying assumptions (1)-(3). Then for any node j of that network the function

(2.5)
$$m_j^{\infty}(t) = \sum_{i \in I} \int_0^{\infty} \lambda_i(t-s) P_{ij}(s) ds$$

is uniform almost periodic and as $t \to \infty$

$$|m_j^{\infty}(t) - m_j(t)| \to 0$$

There exists some constant $B = B(C, a, \delta) > 0$ such that for any network from the class $\mathcal{M}(C, a, \delta)$ the inequality

$$\sup_{t} \left| m_j^{\infty}(t) - d_j \right| \le B\delta^N$$

holds for any node j with dist(j, I) > N.

Proof. We use the following obvious representation, see [1], for the means

$$m_j(t) = \sum_{i \in I} \int_0^t \lambda_i(s) P_{ij}(t-s) ds = \sum_{i \in I} \int_0^t \lambda_i(t-s) P_{ij}(s) ds$$

By (2.2) we have

$$\int_0^\infty P_{ij}(s)ds = \sum_{n=1}^\infty \sum_{L \in \mathbf{L}_n(i,j)} P_n(L) \int_0^\infty ds \int_0^s dF_L(u) (1 - F_j(s - u))$$

Integrating by parts in the above integral we get that for all paths L to the node j

$$\int_0^\infty ds \int_0^s dF_L(u)(1 - F_j(s - u)) = \int_0^\infty s dF_j(s) = \mu_j$$

and, hence,

(2.6)
$$\int_0^\infty P_{ij}(s)ds = \sum_{n=1}^\infty p^{(n)}(i,j)\mu_j = b(i,j)\mu_j$$

Then, as $t \to \infty$,

$$|m_j^{\infty}(t) - m_j(t)| = \left| \sum_{i \in I} \int_t^{\infty} \lambda_i(t - s) P_{ij}(s) ds \right| \to 0,$$

as the functions $\lambda_i(t)$ are bounded and for each network the set of input nodes is finite. Substituting Fourier series (2.3) in (2.5) and using the uniform convergence of the Fourier series we get

(2.7)
$$m_j^{\infty}(t) = \sum_{i \in I} c_0(i) \int_0^{\infty} P_{ij}(s) ds + \sum_{i \in I} \sum_{k=1}^{\infty} c_k(i) \exp(i \sigma_{k,i} t) g_{ij}(\sigma_{k,i}),$$

where

(2.8)
$$g_{ij}(\sigma) = \int_0^\infty \exp(-i\,\sigma s)\,P_{ij}(s)ds,$$

It follows from (2.6) that

(2.9)
$$\sum_{i \in I} c_0(i) \int_0^\infty P_{ij}(s) ds = d_j$$

Let dist(j, I) > N. According to (2.2), we have

$$P_{ij}(s) = \sum_{n>N} \sum_{L \in \mathbf{L}_n(i,j)} P_n(L) G_L(s),$$

where we have denoted

$$G_L(s) = \int_0^s dF_L(u)(1 - F_j(s - u)),$$

Integrating by parts in (2.8) we find

$$g_{ij}(\sigma) = \frac{1}{i\sigma} \int_0^\infty \exp(-i\sigma s) dP_{ij}(s),$$

because $P_{ij}(0) = 0$ and $P_{ij}(s) \to \infty$ as $s \to \infty$. Note that

$$dP_{ij}(s) = \sum_{n>N} \sum_{L \in \mathbf{L}_n(i,j)} P_n(L) dG_L(s)$$

and $dG_L(s) = dF_L(s) - (dF_L \star dF_j)(s)$. Therefore,

$$g_{ij}(\sigma) = \sum_{n>N} \sum_{L \in \mathbf{L}_n(i,j)} P_n(L) \left(\frac{\psi_L(-\sigma) - \psi_L(-\sigma) \psi_j(-\sigma)}{i\sigma} \right),$$

where

(2.10)
$$\psi_L(-\sigma) = \psi_{i_1, i_2}(-\sigma) \dots \psi_{i_{n-1}, i_n}(-\sigma),$$

and is ψ_i the characteristic function of distribution F_i . Finally, we get

(2.11)
$$g_{ij}(\sigma) = \sum_{n>N} \sum_{L \in \mathbf{L}_n(i,j)} P(L)\psi_L(-\sigma) \left(\frac{1 - \psi_j(-\sigma)}{i\sigma}\right).$$

Further, by **conditions** (ii) and (iii) we have that for all $\sigma_{k,i}$

$$|\psi_L(-\sigma_{k,i})| \leq \delta^{|L|}$$

where |L| is the length of the path L and, thus, for all i, j, k with dist(j, I) > N we have

$$(2.12) |g_{ij}(\sigma_{k,i})| \le \frac{2}{a} \sum_{n>N} p^{(n)}(i,j)\delta^n \le \frac{2}{a(1-\delta)}\delta^N$$

Hence, using formulas (2.7),(2.9) and **condition** (i) we have

$$\left| m_j^{\infty}(t) - d_j \right| \le \sum_{i \in I} \sum_{k=1}^{\infty} |c_k(i)| \frac{2}{a(1-\delta)} \delta^N = B(C, a.\delta) \delta^N$$

The theorem is proved.

Remark 1. Our model can be considered as an open queueing network with M nodes and infinite number of services available at each node. Different questions for these networks were studied in [1, 2, 13].

However, we should note that the property of almost periodicity plays a key role in theorem 1. It is not difficult to give examples of non almost periodic functions $\lambda_i(t)$ for which the theorem does not hold.

3. Stochastic stationary input

Here we assume that the input rates form a random stationary vector process (random media)

$$(\lambda_1(t,\omega_{env}),...,\lambda_{|I|}(t,\omega_{env}))$$

The components are assumed to be mutually independent, and almost surely $|(\lambda_i(t, \omega_{env}))| \le L_i$ for some constants L_i . Then the mean numbers of molecules depend on ω_{env} , and we will study their typical behavior.

We will denote E_{env} the mean over the random environment. Let $\lambda_i = E_{env}\lambda_i(t)$ and $r_i(t)$ be the correlation function: $r_i(t) = E_{env}(\lambda_i(t+s) - \lambda_i)(\lambda_i(s) - \lambda_i), i = 1, ..., N$. We assume that $r_i(t)$ is continuous at 0. Let $z_i(d\sigma)$ be the corresponding spectral measure, i.e.

$$r_i(t) = \int_{-\infty}^{+\infty} e^{it\sigma} z_i(d\sigma)$$

We will need the following spectral representation

$$\lambda_i(t) = \lambda_i + \int_{-\infty}^{+\infty} e^{it\sigma} \varsigma_i(d\sigma),$$

where $\varsigma_i(d\sigma)$ is the orthogonal random measure such that $E_{env}\varsigma_i(A) = 0$, $E_{env}|\varsigma_i(A)|^2 = 0$ $z_i(A)$ for any Borel subset $A \subset (-\infty, +\infty)$.

Introduce the random process $l_i(t) = l_i(t, \omega_{env})$ which is the mean number of molecules at node j under a given ω_{env} :

$$l_j(t) = \sum_{i \in I} \int_0^t \lambda_i(t-s) P_{ij}(s) ds,$$

where $P_{ij}(s)$ is defined by formula (2.2). Introduce also the stationary random processes

(3.1)
$$l_j^{\infty}(t) = \sum_{i \in I} \int_0^{\infty} \lambda_i(t-s) P_{ij}(s) ds$$

We define the class $\mathcal{M}_s(C, a, \delta)$, $C, a > 0, 0 < \delta < 1$ consisting of networks satisfying the above condition (iii) and the following conditions

(iv)

$$\operatorname{Var}\left(\sum_{i\in I}\lambda_i(t)\right) \leq C$$

or, equivalently, in terms of the spectral measure $\sum_{i \in I} \int_{-\infty}^{+\infty} z_i(d\sigma) \leq C$. (v) there exists spectral gap, i. e. for all i the spectral measure $z_i(A_0) = 0$ for some interval $A_0 = (-a, a)$.

We define also another class $\mathcal{M}_{ss}(C, a, \delta), C, a > 0, 0 < \delta < 1$, consisting of networks satisfying the above conditions (iii),(iv) and also the following condition

(vi) for any $\epsilon > 0$ there is a > 0 (small enough) such that

$$\sum_{i \in I} \int_{-a}^{+a} z_i(d\sigma) \le \epsilon$$

and b(i,j) and μ_i are uniformly bounded: $b(i,j) \leq b, \mu_i \leq \mu$ for some constants $b, \mu > 0$. It follows from this condition that the spectral measures $z_i(d\sigma)$ have no atoms at 0.

Theorem 2. Assume that conditions 2–3 hold. Then, for any j, as $t \to \infty$

$$l_i^{\infty}(t) - l_j(t) \to 0$$
 with probability 1.

There exists the constant $\hat{B} = \hat{B}(C, a, \delta) > 0$ such that for any network from the class $\mathcal{M}_s(C, a, \delta)$ the inequality

$$E_{env} \left| l_j^{\infty}(t) - e_j \right|^2 \le \hat{B} \delta^{2N}, \qquad e_j = E_{env} l_j^{\infty}(t)$$

holds for any node j with dist(j, I) > N.

For any $\epsilon > 0$ there exists N_0 large enough such that for any network from the class $\mathcal{M}_{ss}(C, a, \delta)$ the inequality

$$E_{env} \left| l_j^{\infty}(t) - e_j \right|^2 \le \epsilon$$

holds for any node j with dist(j, I) > N.

Proof. To prove the first assertion of the theorem let us note that

$$|l_j^{\infty}(t) - l_j(t)| = \left| \sum_{i \in I} \int_t^{\infty} \lambda_i(t - s) P_{ij}(s) ds \right| \le \sum_{i \in I} L_i \int_t^{\infty} P_{ij}(s) ds \to 0,$$

as the sum is finite and the integrals are convergent.

Prove the second assertion of the theorem. Using (3.1) and (2.11) we get the following spectral representation for the process $l_i^{\infty}(t)$

$$l_j^{\infty}(t) = E_{env}l_j^{\infty}(t) + \sum_{i \in I} \int_{-\infty}^{+\infty} e^{it\sigma} g_{ij}(\sigma) \varsigma_i(d\sigma),$$

where

$$E_{env}l_j^{\infty}(t) = \sum_{i \in I} \lambda_i \int_0^{\infty} P_{ij}(s)ds = \sum_{i \in I} \lambda_i b(i,j)\mu_j$$

and $g_{ij}(\sigma)$ is defined by formula (2.8). Hence, $\sum_{i \in I} g_{ij}^2(\sigma) z_i(d\sigma)$ is the spectral measure of $l_i^{\infty}(t)$, and, thus,

$$D_j = E_{env} \left| l_j^{\infty}(t) - E_{env} l_j^{\infty}(t) \right|^2 = \sum_{i \in I} \int_{-\infty}^{+\infty} g_{ij}^2(\sigma) z_i(d\sigma),$$

Let dist(j, I) > N. By formula (2.12) and condition (iv) we have

(3.2)
$$\sum_{i \in I} \int_{|\sigma| > a} g_{ij}^2(\sigma) z_i(d\sigma) \le \hat{B}(C, a, \delta) \delta^{2N}$$

By condition (v) we get

$$D_j = E_{env} \left| l_j^{\infty}(t) - E_{env} l_j^{\infty}(t) \right|^2 \le \hat{B}(C, a, \delta) \delta^{2N}$$

To prove the third assertion note that for $\epsilon > 0$ arbitrary small we can choose $\kappa > 0$ small enough such that

$$\sum_{i \in I} \int_{-\kappa}^{+\kappa} g_{ij}^2(\sigma) z_i(d\sigma) \le \epsilon$$

This follows from the condition (vi). If $dist(j,I) > N_0$, then, as follows from reasoning similar to (3.2), the expression $\sum_{i \in I} \int_{|\sigma| > \kappa} g_{ij}^2(\sigma) z_i(d\sigma)$ can be made arbitrary small for N_0 large enough.

The theorem is proved.

4. Two thermodynamic limits

There are two quite different thermodynamic limits in our context. The first one is quite common: instead of considering finite networks we could consider infinite networks.

The second one corresponded to the fact that in chemical kinetics it is normally assumed that the number of molecules in a given compartment is large. More exactly, we consider a family $\xi^{(\Lambda)}$ of the processes defined above with the same graph G, parametrized by the large parameter Λ interpreted as the volume of the compartment. For simplicity of notation we assume that all compartments have the same volume.

This parameter appears in our model in two places. Firstly, we assume that the initial conditions for the means $m_{\alpha,v}^{(\Lambda)}(0)$ are such that at time zero there is a limit

$$\lim_{\Lambda \to \infty} \Lambda^{-1} m_{\alpha,v}^{(\Lambda)}(0) = c_{\alpha,v}(0)$$

(initial concentrations). Secondly, the input rates for the process $\xi^{(\Lambda)}$ are scaled as

$$\lambda_i(t) = a_i(t)\Lambda$$

for some non negative functions $a_i(t)$,

As the means satisfy the following equations (for the Markov case)

$$\frac{dm_i(t)}{dt} = \lambda_i(t) + \sum_i (m_j(t)\lambda_{ji} - m_i(t)\lambda_{ij}) - m_i(t)\lambda_{i0}$$

the limiting concentrations, due to linearity,

$$c_i(t) = \lim_{\alpha \to 0} \Lambda^{-1} m_{\alpha,\nu}^{(\Lambda)}(t), \quad i = (\alpha, \nu)$$

exist and satisfy the same equations, see below. Note that these equations coincide with the classical chemical kinetics equations for the system with uni modular reactions.

One can take into account the possibility that if a molecule goes from compartment α to compartment β , it may be absent for some time from both compartments. This case can be reduced to the previous case if we introduce a new vertex v_{ij} for any edge (i, j) with nonzero λ_{ij} .

Instead of considering the case with general F_{ij} , it is more reasonable to modify the Markov model slightly. First of all we assume that when a particle chooses transition to a different compartment it does not change type, that is transitions $(\alpha, v) \to (\beta, w)$, $\alpha \neq \beta$, are only possible when v = w. Moreover, it appears in compartment β only after random time τ_{ik} , $i = (\alpha, v)$, $k = (\beta, v)$. All times τ_{ik} are independent and have distribution functions F_{ik} . Then the limiting equations are

$$\frac{dc_k(t)}{dt} = l_k(t) + \sum_{\substack{i = (\alpha, v) : \\ \alpha = \beta}} c_i(t)\lambda_{ik} + \sum_{\substack{i = (\alpha, v) : \\ \alpha \neq \beta}} \int_0^t c_i(t - u)\lambda_{ik}dF_{ik}(u) - \sum_{i = (\alpha, v) : \alpha \neq \beta} c_i(t)\lambda_{kj} - c_k(t)\lambda_{k0}, \qquad k = (\beta, w).$$

For these cumbersome chemical kinetics equations, we can prove the following result.

Corollary 4.1. Under conditions of theorem 1 there exist almost periodic functions $c_j^{\infty}(t)$, j = 1, ..., J, and constants d_j , such that as $t \to \infty$

$$|c_j(t) - c_i^{\infty}(t)| \to 0$$
 and $\sup |c_i^{\infty}(t) - d_j| \le B\delta^N$

for j satisfying the condition dist(j, N) > N.

Remark 2. Chemical networks were considered in many papers, some models close to ours one can find in [5, 8, 6, 7]. A microscopic model of one thermodynamic compartment was considered in [3]. Network of compartments were introduced in [4]. It is important to extend our results to more general networks. Using deep results from [11, 12] one could possibly consider networks with FIFO queues. Most important however is to study binary reactions case.

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